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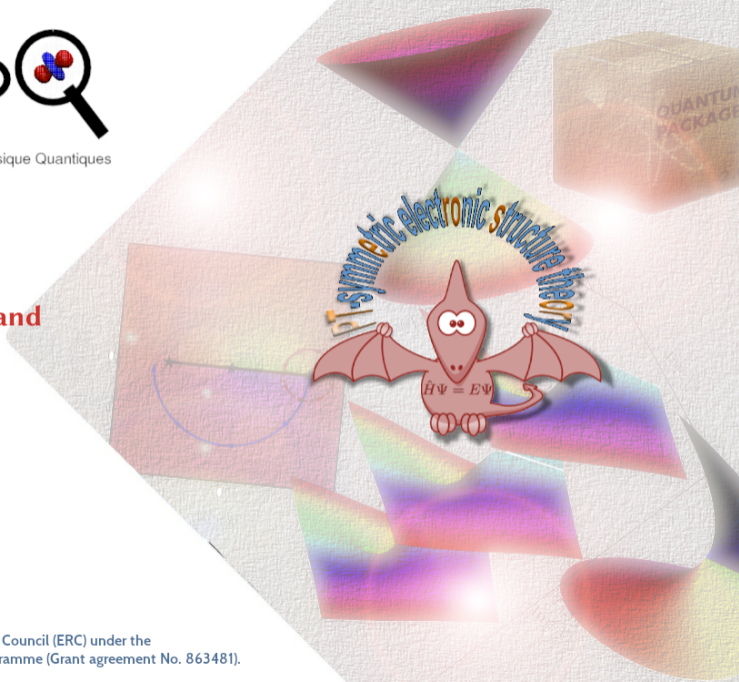
CIPSI: selected configuration interaction methods for ground and excited states

Pierre-François Loos & Friends

19th September 2023

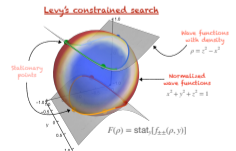
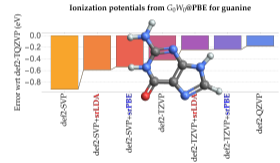
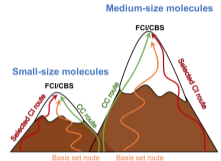
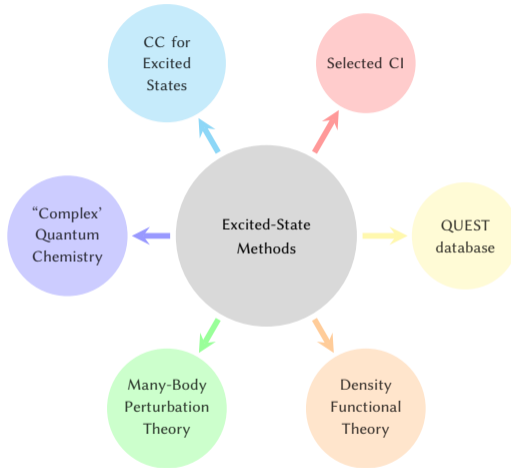
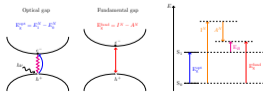
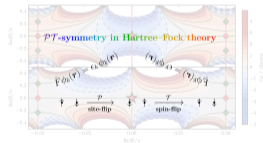
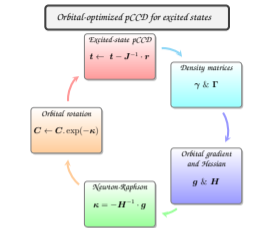
Laboratoire de Chimie et Physique Quantiques, IRSAMC, UPS/CNRS, Toulouse

<https://lcpq.github.io/pterosor>



PTEROSOR has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (Grant agreement No. 863481).

General Overview of our Research Group



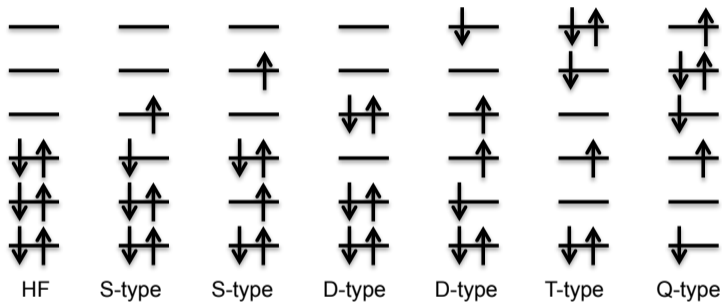
<https://lcpq.github.io/PTEROSOR/>

- ▶ This is the **oldest** and perhaps the **easiest method to understand**
- ▶ The CI wave function is a **linear combination of Slater determinants**
- ▶ CI methods use **excited determinants** to “improve” the reference (usually Hartree-Fock) wave function

$$|\Phi_{\text{CI}}\rangle = \underbrace{c_0 |\Psi_0\rangle}_{\text{reference}} + \underbrace{\sum_i^a c_i^a |\Psi_i^a\rangle}_{\text{singles}} + \underbrace{\sum_{\substack{i<j \\ a<b}} c_{ij}^{ab} |\Psi_{ij}^{ab}\rangle}_{\text{doubles}} + \underbrace{\sum_{\substack{i<j<k \\ a<b<c}} c_{ijk}^{abc} |\Psi_{ijk}^{abc}\rangle}_{\text{triples}} + \underbrace{\sum_{\substack{i<j<k<l \\ a<b<c<d}} c_{ijkl}^{abcd} |\Psi_{ijkl}^{abcd}\rangle}_{\text{quadruples}} + \dots$$

- ▶ CI is based on the **variational principle** (like the Hartree-Fock approximation)

Excited determinants



CI wave function

$$|\Phi_{CI}\rangle = c_0 |0\rangle + c_S |S\rangle + c_D |D\rangle + c_T |T\rangle + c_Q |Q\rangle + \dots$$

- ▶ When $|S\rangle$ (**singles**) are taken into account: **CIS**

$$|\Phi_{\text{CIS}}\rangle = c_0 |0\rangle + c_S |S\rangle$$

NB: CIS is an **excited state method**

- ▶ When $|S\rangle$ and $|D\rangle$ are taken into account: **CISD**

$$|\Phi_{\text{CISD}}\rangle = c_0 |0\rangle + c_S |S\rangle + c_D |D\rangle$$

NB: CISD is the **most commonly-used** CI method

- ▶ When $|S\rangle$, $|D\rangle$ and $|T\rangle$ (**triples**) are taken into account: **CISDT**

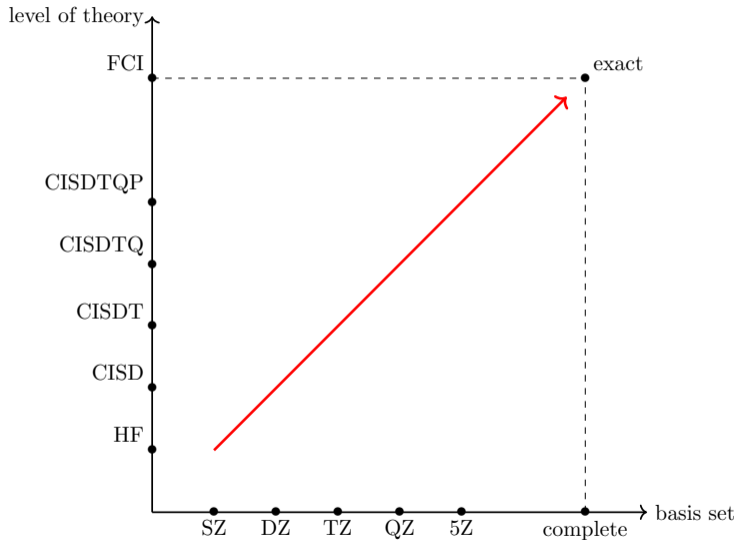
$$|\Phi_{\text{CISDT}}\rangle = c_0 |0\rangle + c_S |S\rangle + c_D |D\rangle + c_T |T\rangle$$

- ▶ **CISDTQ**, etc.

- ▶ When all possible excitations are taken into account, **this is called a Full CI calculation (FCI)**

$$|\Phi_{\text{FCI}}\rangle = c_0 |0\rangle + c_S |S\rangle + c_D |D\rangle + c_T |T\rangle + c_Q |Q\rangle + \dots$$

- ▶ FCI gives the **exact solution of the Schrödinger equation within a given basis**
- ▶ FCI is becoming more and more fashionable these days (e.g. FCIQMC and CIPSI methods)
- ▶ So, why do we care about other methods?
- ▶ **Because FCI is super computationally expensive!**



“Assume we have 10 electrons in 38 spin MOs: 10 are occupied and 28 are empty”

- ▶ There is C_{10}^k possible ways of selecting k electrons out of the 10 occupied orbitals

$$C_n^k = \frac{n!}{k!(n-k)!}$$

- ▶ There is C_{28}^k ways of distributing them out in the 28 virtual orbitals
- ▶ For a given excitation level k , there is $C_{10}^k C_{28}^k$ excited determinants
- ▶ The total number of possible excited determinant is

$$\sum_{k=0}^{10} C_{10}^k C_{28}^k = C_{38}^{10} = 472,733,756$$

- ▶ This is a lot...

For $N = 10$ and $K = 38$:

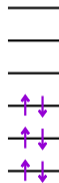
k	Num. of excitations
0	1
1	280
2	17,010
3	393,120
4	4,299,750
5	24,766,560
6	79,115,400
7	142,084,800
8	139,864,725
9	69,069,000
10	13,123,110
Tot.	472,733,756

How to “span” the Hilbert space: Excitation-based CI

e	
0	
1	
2	
3	

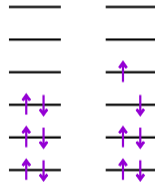
How to “span” the Hilbert space: Excitation-based CI

e	
0	HF
1	
2	
3	



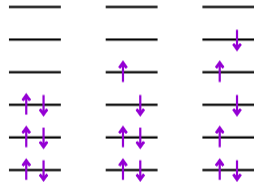
How to “span” the Hilbert space: Excitation-based CI

e	
0	
1	CIS
2	
3	



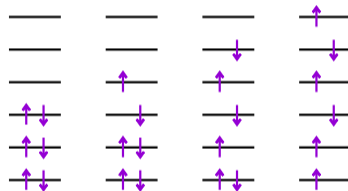
How to “span” the Hilbert space: Excitation-based CI

e	
0	
1	
2	CISD
3	



How to “span” the Hilbert space: Excitation-based CI

e	
0	
1	
2	
3	CISDT

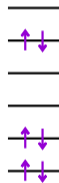


How to “span” the Hilbert space: Seniority-based CI

s	0	2	4	6

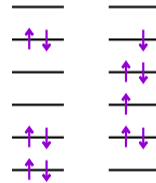
How to “span” the Hilbert space: Seniority-based CI

s	0	2	4	6
	sCI0			



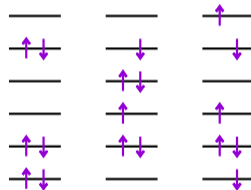
How to “span” the Hilbert space: Seniority-based CI

s	0	2	4	6
		sCI2		



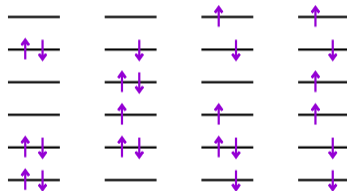
How to “span” the Hilbert space: Seniority-based CI

s	0	2	4	6
			sCI4	



How to “span” the Hilbert space: Seniority-based CI

s	0	2	4	6
				sCI6



e/s	0	2	4	6	8
0					
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0	HF				
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1		CIS			
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2			CISD		
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2					
3				CISDT	
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0	sCI0				
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1		sCI2			
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2			sCI4		
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2					
3				sCI6	
4					
5					
6					

Hierarchy CI (hCI)

$$h = \frac{e + s/2}{2}$$

- ▶ e : excitation degree
- ▶ s : seniority number
- ▶ h : hierarchy parameter

e/s	0	2	4	6	8
0					
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0	HF				
1					
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1		hCI1			
2					
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2		hCI1.5			
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2			hCl2		
3					
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2					
3			hCl2.5		
4					
5					
6					

e/s	0	2	4	6	8
0					
1					
2					
3				hCl3	
4					
5					
6					

Excitation-based CI vs Hierarchy CI vs Seniority-based CI

e/s	0	2	4	6
0	HF			
1		CIS		
2			CISD	
3				CISDT
4				
5				
6				

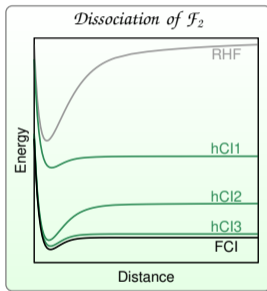
e/s	0	2	4	6
0	HF			
1		hCI1		
2		hCI1.5	hCI2	
3			hCI2.5	hCI3
4				
5				
6				

e/s	0	2	4	6
0				
1				
2				
3				
4				
5				
6	sCI0	sCI2	sCI4	sCI6

Hierarchy configuration interaction (hCI)

Excitation degree e
Seniority number s
Hierarchy parameter $h = \frac{e+s/2}{2}$

e/s	0	2	4	6
0	HF			
1		hCI1		
2		hCI1.5	hCI2	
3			hCI2.5	hCI3
4				
5				
6				



Fábri Kossoski

Kossoski, Damour & Loos, JPLC 13 (2022) 4342 (see also arXiv:2308.14618)

Selected Configuration Interaction: “sparse” exploration of the FCI space

- *“Among the very large number of determinants contained in the FCI space, only a tiny fraction of them significantly contributes to the energy”*

CIPSI = CI using a Perturbative Selection made Iteratively

- ▶ Developed in Toulouse many (many) years ago
Huron, Malrieu & Rancurel, JCP 58 (1973) 5745
- ▶ Based on old idea by Bender and Davidson, and Whitten and Hackmeyer
Bender & Davidson, Phys. Rev. 183 (1969) 23
Whitten & Hackmeyer, JCP 51 (1969) 5584
- ▶ CIPSI (and SCI methods in general) has been recently resurrected!
Giner, Scemama & Caffarel, CJC 91 (2013) 879
Giner, Scemama & Caffarel, JCP 142 (2015) 044115
- ▶ CIPSI \approx heat-bath CI (Umrigar) \approx adaptive sampling CI (Evangelista) \approx iterative CI (Liu) \approx incremental CI (Zimmerman) \approx FCIQMC (Alavi)

Volume 126, number 1

CHEMICAL PHYSICS LETTERS

25 April 1986

**CONVERGENCE OF A MULTIREFERENCE SECOND-ORDER MBPT METHOD (CIPSI)
USING A ZERO-ORDER WAVEFUNCTION DERIVED FROM AN MS SCF CALCULATION** ☆

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Received 28 November 1985; in final form 7 February 1986

Approximate natural orbitals and the convergence of a second order multireference many-body perturbation theory (CIPSI) algorithm^{a)}

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Departamento Química, Facultat de Química de Tarragona, Universitat de Barcelona, Pça. Imperial Tàrraco 1, 43005-Tarrogona, Spain

(Received 6 June 1988; accepted 9 August 1988)

Theor Chim Acta (1992) 82: 229–238

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Chimica Acta**

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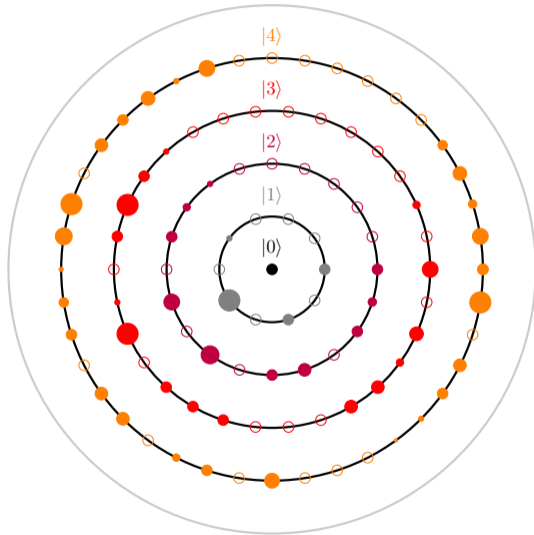
Treating large intermediate spaces in the CIPSI method through a direct selected CI algorithm

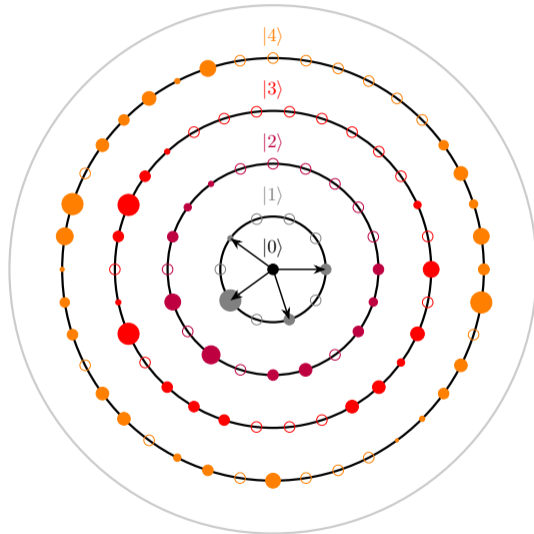
A. Povill, J. Rubio, and F. Illas

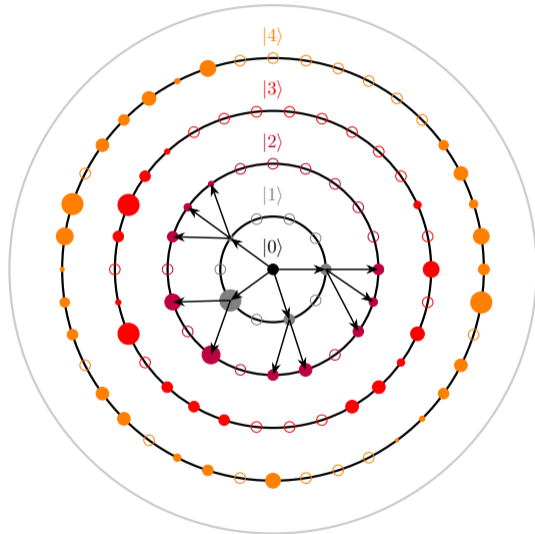
Departament de Química Física, Grup de Química Quàntica, Universitat de Barcelona,
C/Martí i Franquès 1, E-08028 Barcelona, Spain

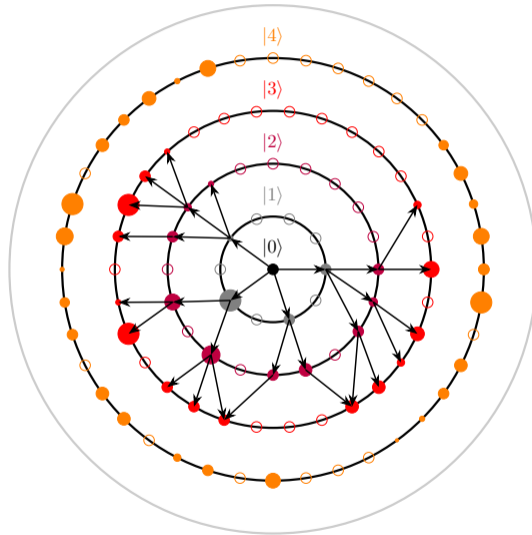
Table 1. Summary of results for HF including the dimension of the $\{G\}$, $\{M\}$ and $\{GD\}$ subspaces, the variational MRCI energy, E_{MRCI} , and the error with respect to the FCI energy, ΔE . Energies are in atomic units

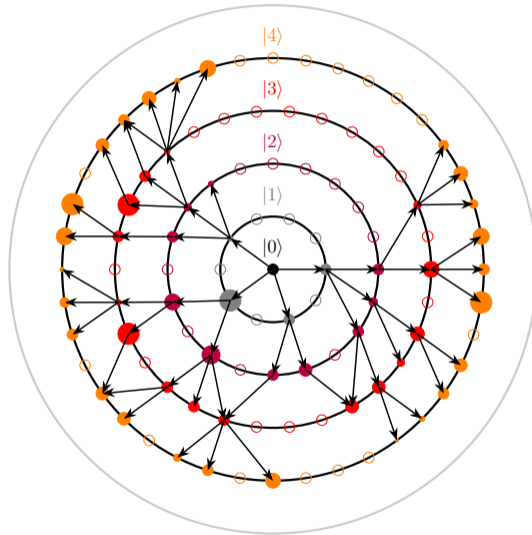
dim $\{G\}$	dim $\{GD\}$	E_{MRCI}	ΔE
113	81868	-100.272466	0.001641
194	125016	-100.272912	0.001195
278	159009	-100.273067	0.001040
388	191984	-100.273068	0.001039
FCI [28]	1.3×10^7	-100.274107	—











- ▶ **Green:** reference/variational/internal wave function (zeroth-order or model space)
- ▶ **Red:** perturbations or external wave function (first-order or perturbative space)

$$\begin{array}{l} \text{internal space } |I\rangle \\ \text{perturbations } |\alpha\rangle \end{array} \left\{ \begin{array}{|c|c|} \hline \mathbf{H}^{(0)} & \mathbf{h}^\dagger \\ \hline \mathbf{h} & \mathbf{H}^{(1)} \\ \hline \end{array} \right.$$

1. Define a (zeroth-order) *reference* wave function:

$$|\Psi^{(0)}\rangle = \sum_{I \in \mathcal{I}} c_I |I\rangle \qquad E^{(0)} = \frac{\langle \Psi^{(0)} | \hat{H} | \Psi^{(0)} \rangle}{\langle \Psi^{(0)} | \Psi^{(0)} \rangle}$$

2. Generate *external determinants*:

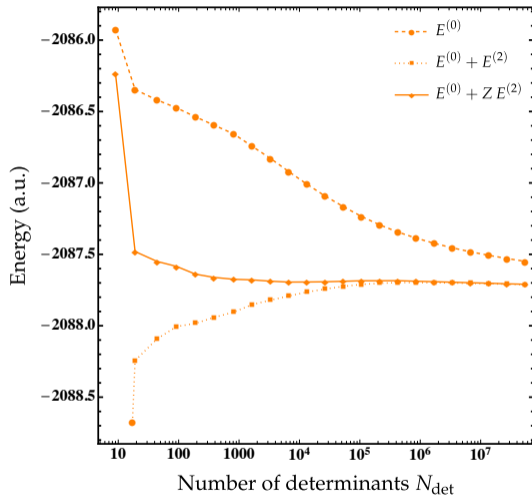
$$\mathcal{A} = \{(\forall I \in \mathcal{I}) (\forall \hat{T} \in \mathcal{T}_1 \cup \mathcal{T}_2) : |\alpha\rangle = \hat{T}|I\rangle\}$$

3. Second-order perturbative contribution of each $|\alpha\rangle$:

$$\delta E(\alpha) = \frac{|\langle \Psi^{(0)} | \hat{H} | \alpha \rangle|^2}{E^{(0)} - \langle \alpha | \hat{H} | \alpha \rangle}$$

4. Select $|\alpha\rangle$ with largest $\delta E(\alpha)$ and add them to \mathcal{I}
5. Diagonalize \hat{H} in $\mathcal{I} \Rightarrow$ update $|\Psi^{(0)}\rangle$ and $E^{(0)}$
6. Iterate

Ground state of Cr_2 in cc-pVQZ: full-valence CAS(28e,198o)



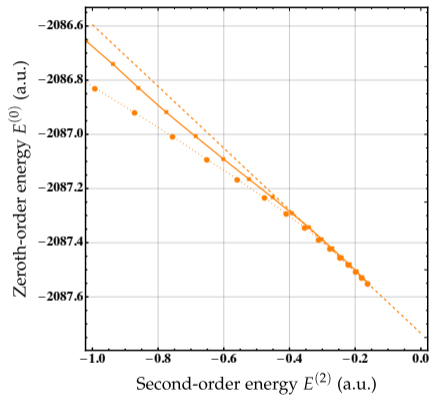
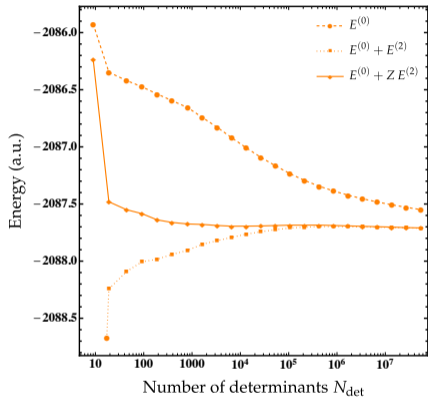
Garniron et al., JCTC 15 (2019) 3591

- ▶ Second-order Epstein-Nesbet correction:

$$E^{(2)} = \sum_{\alpha} \delta E(\alpha)$$

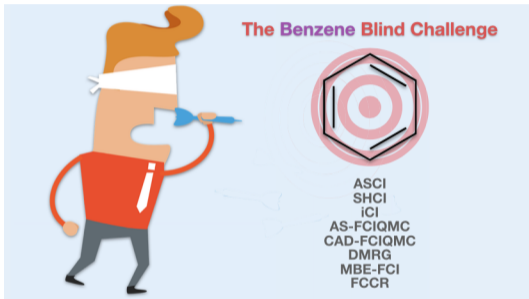
- ▶ $|\alpha\rangle$'s with largest $\delta E(\alpha)$ have been added to $\Psi^{(0)}$ previously
⇒ only small contributions remaining
- ▶ \mathcal{A} increases with \mathcal{I}
⇒ a *very large* number of *very small* contributions
- ▶ In practice, we use a semi-stochastic algorithm to compute $E^{(2)}$
⇒ *much faster!!*
Garniron, Scemama, Loos & Caffarel, JCP 147 (2017) 034101
- ▶ We *linearly extrapolate* to $E^{(2)} = 0$ to reach the FCI limit (exFCI)

Ground state of Cr_2 in cc-pVQZ: full-valence CAS(28e,198o)

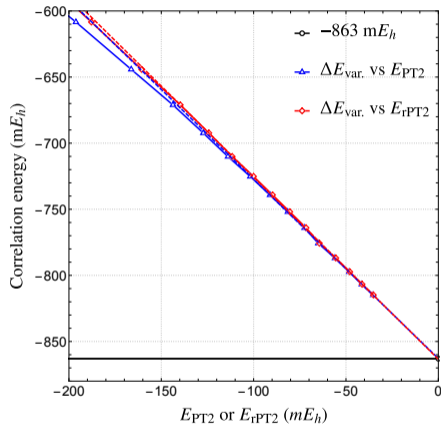
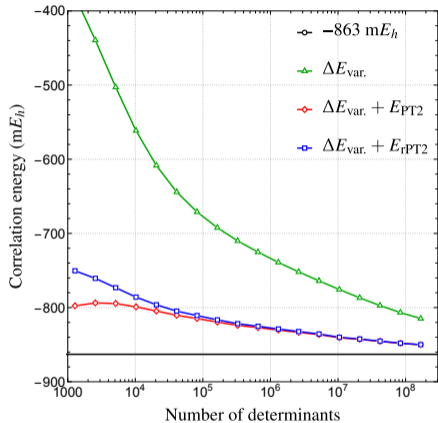


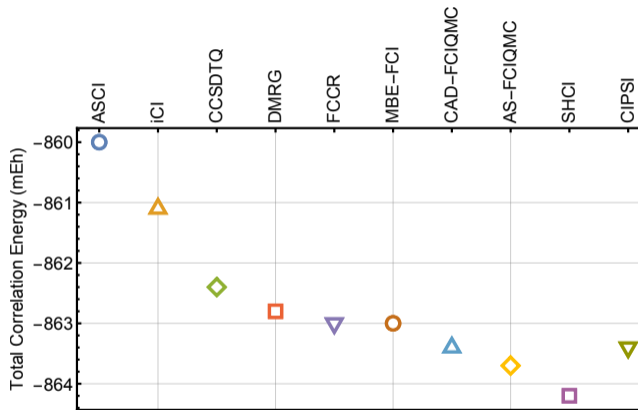
Garniron et al., JCTC 15 (2019) 3591

The Benzene Blind Challenge: Frozen-core correlation energy (cc-pVDZ)



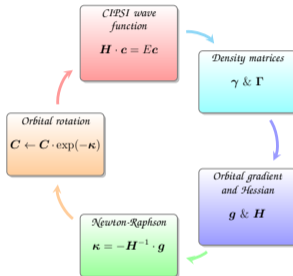
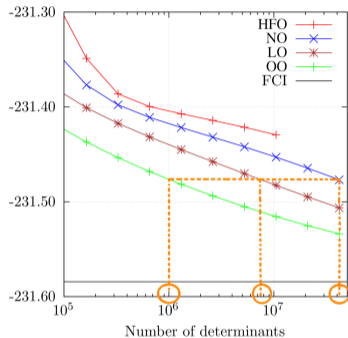
Eriksen et al. JPCL 11 (2020) 8922





Loos, Damour & Scemama, JCP 153 (2020) 176101

Orbital-optimized CIPSI for $C_6H_6/cc\text{-pVDZ}$ (and many others)



- ▶ Orbital optimization largely accelerates the convergence of selected CI
- ▶ Trust-region Newton-Raphson algorithm

Damour, V ril, Kossoski, Caffarel, Jacquemin, Scemama & Loos, JCP 155 (2020) 176101



Yann Damour

Toward a systematic improvement of the fixed-node approximation in diffusion Monte Carlo for solids—A case study in diamond

Cite as: J. Chem. Phys. **153**, 184111 (2020); <https://doi.org/10.1063/5.0021036>

Submitted: 06 July 2020 . Accepted: 12 October 2020 . Published Online: 11 November 2020

 Anouar Benali,  Kevin Gasperich,  Kenneth D. Jordan, Thomas Applencourt,  Ye Luo,  M. Chandler Bennett,  Jaron T. Krogel,  Luke Shulenburger,  Paul R. C. Kent,  Pierre-François Loos,  Anthony Scemama, and  Michel Caffarel

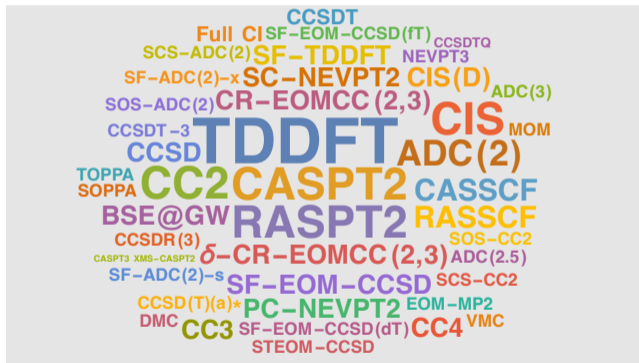


Michel Caffarel

Zoo of functionals...



And this is just for excited states...



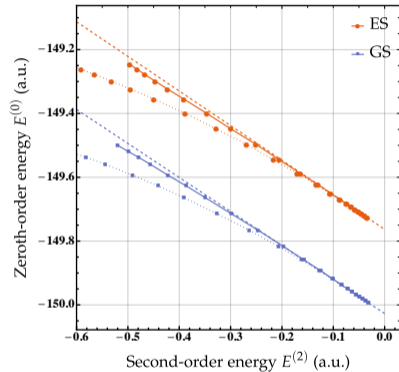
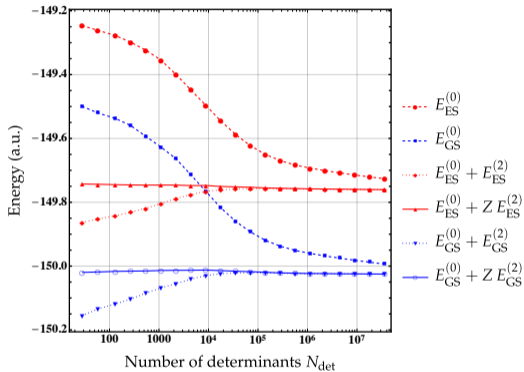


Table 1. Zeroth-Order Energy $E^{(0)}$, Second-Order Perturbative Correction $E^{(2)}$, and Its Renormalized Version $\text{ZE}^{(2)}$ (in hartree) of CN3 for Increasingly Large Wave Functions^a

N_{det}	$E^{(0)}$		$E^{(0)} + E^{(2)}$			$E^{(0)} + \text{ZE}^{(2)}$		
	GS (a.u.)	ES (a.u.)	GS (a.u.)	ES (a.u.)	$\Delta E(\text{eV})$	GS (a.u.)	ES (a.u.)	$\Delta E(\text{eV})$
28	-149.499574	-149.246268	-150.155(1)	-149.863(1)	7.95(5)	-150.020(1)	-149.743(1)	7.54(5)
58	-149.519908	-149.261390	-150.134(1)	-149.853(1)	7.67(5)	-150.018(1)	-149.744(1)	7.48(5)
131	-149.537424	-149.277496	-150.118(1)	-149.8427(9)	7.52(4)	-150.017(1)	-149.7449(9)	7.39(4)
268	-149.559465	-149.298484	-150.1035(9)	-149.8308(9)	7.42(4)	-150.0158(9)	-149.7457(9)	7.35(4)
541	-149.593434	-149.323302	-150.0845(8)	-149.8186(8)	7.24(4)	-150.0152(8)	-149.7463(8)	7.32(4)
1101	-149.627202	-149.354807	-150.0683(8)	-149.8045(8)	7.18(3)	-150.0137(8)	-149.7460(8)	7.28(3)
2207	-149.663850	-149.399522	-150.0549(7)	-149.7879(7)	7.26(3)	-150.0132(7)	-149.7462(7)	7.27(3)
4417	-149.714222	-149.448133	-150.0409(6)	-149.7762(6)	7.20(3)	-150.0130(6)	-149.7478(6)	7.22(3)
8838	-149.765886	-149.496401	-150.0296(5)	-149.7655(5)	7.19(2)	-150.0124(5)	-149.7473(5)	7.21(2)
17 680	-149.817301	-149.545048	-150.0239(4)	-149.7615(4)	7.14(2)	-150.0141(4)	-149.7505(4)	7.17(2)
35 380	-149.859737	-149.587668	-150.0216(3)	-149.7582(3)	7.17(1)	-150.0161(3)	-149.7518(3)	7.19(1)
70 764	-149.893273	-149.623235	-150.0207(2)	-149.7566(3)	7.18(1)	-150.0174(2)	-149.7530(3)	7.19(1)
141 545	-149.919463	-149.650109	-150.0214(2)	-149.7572(2)	7.189(8)	-150.0194(2)	-149.7550(2)	7.196(8)
283 108	-149.937839	-149.669735	-150.0224(2)	-149.7576(2)	7.206(7)	-150.0211(2)	-149.7562(2)	7.209(7)
566 226	-149.950918	-149.683278	-150.0233(1)	-149.7580(1)	7.217(6)	-150.0223(1)	-149.7570(1)	7.219(6)
1 132 520	-149.960276	-149.693053	-150.0238(1)	-149.7588(1)	7.212(5)	-150.0231(1)	-149.7580(1)	7.214(5)
2 264 948	-149.968203	-149.700907	-150.0240(1)	-149.7590(1)	7.211(4)	-150.0235(1)	-149.7584(1)	7.212(4)
4 529 574	-149.975230	-149.708061	-150.0245(1)	-149.7594(1)	7.215(4)	-150.0241(1)	-149.7589(1)	7.216(4)
9 057 914	-149.981770	-149.714526	-150.02463(9)	-149.75981(8)	7.206(3)	-150.02434(9)	-149.75948(8)	7.207(3)
18 110 742	-149.987928	-149.720648	-150.02495(7)	-149.76025(8)	7.203(3)	-150.02474(7)	-149.76000(8)	7.204(3)
36 146 730	-149.993593	-149.726253	-150.02527(6)	-149.76065(7)	7.198(3)	-150.02502(6)	-149.760 47(7)	7.198(3)

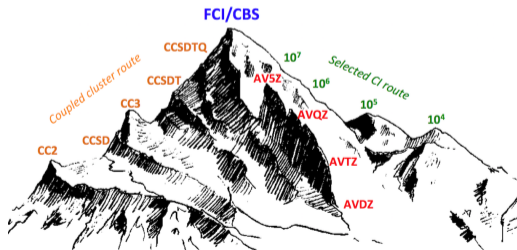
^aThe excitation energy ΔE (in eV) is the energy difference between the ground state (GS) and the excited state (ES). The statistical error, corresponding to one standard deviation, is reported in parentheses.

A Mountaineering Strategy to Excited States: Highly Accurate Reference Energies and Benchmarks

Pierre-François Loos,^{*,†} Anthony Scemama,[†] Aymeric Blondel,[‡] Yann Garniron,[†] Michel Caffarel,[†] and Denis Jacquemin^{*,‡}

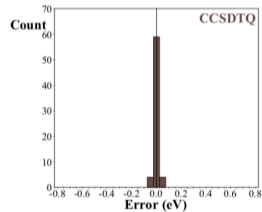
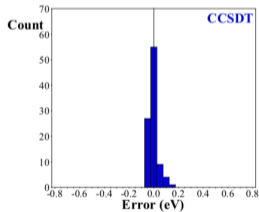
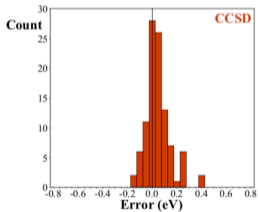
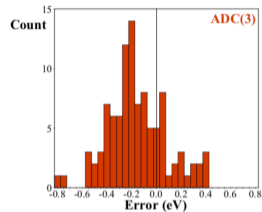
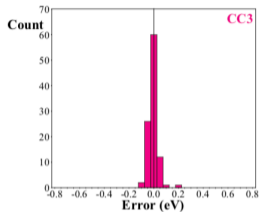
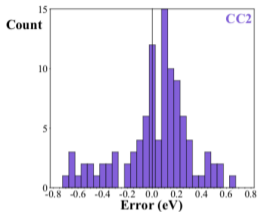
[†]Laboratoire de Chimie et Physique Quantiques, Université de Toulouse, CNRS, UPS, 31013 Toulouse Cedex 6, France

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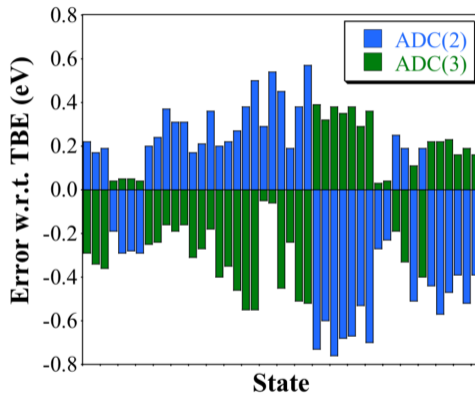


- ▶ 110 **vertical excitation energies** (VTEs) and oscillator strengths
- ▶ 18 small molecules from 1 to 3 non-H atoms
- ▶ CC3/aug-cc-pVTZ geometries
- ▶ mostly singly-excited states and very few doubly-excited states
- ▶ rely on FCI to define **“theoretical best estimates”** (TBEs)
- ▶ aug-cc-pVTZ and CBS vertical energies
- ▶ benchmark popular excited-state methods accounting for double and triple excitations

Benchmarking excited-state methods vs TBE/aug-cc-pVTZ



Errors in ADC(2) & ADC(3) for states with large (> 0.15 eV) ADC(2) error



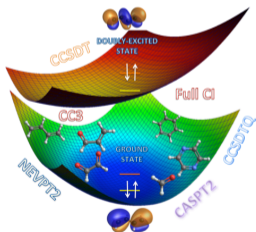
See also Loos & Jacquemin, JPCL 11 (2020) 974

Reference Energies for Double Excitations

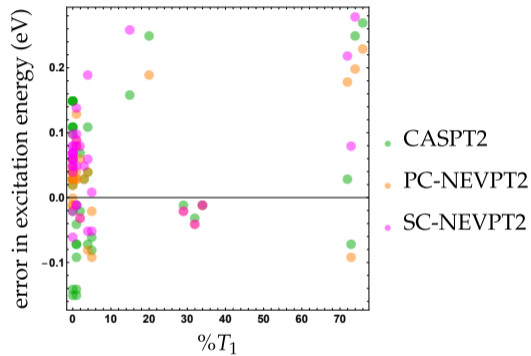
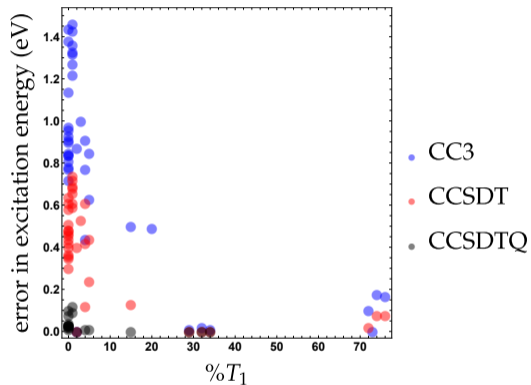
Pierre-François Loos,^{*,†} Martial Boggio-Pasqua,[†] Anthony Scemama,[†] Michel Caffarel,[†] and Denis Jacquemin[‡]

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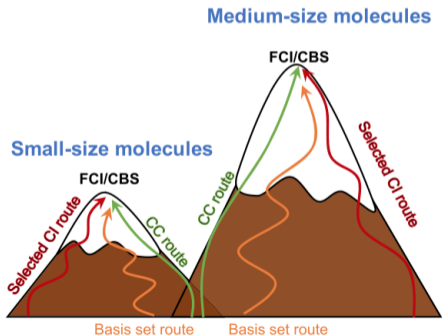


- ▶ 20 VTEs for **doubly-excited states**
- ▶ 14 small- and medium-sized molecules
- ▶ mostly rely on FCI to define TBEs (except for the largest molecules)
- ▶ aug-cc-pVTZ and CBS vertical energies
- ▶ benchmark excited-state methods including **at least** triple excitations
- ▶ additional benchmarks of multi-configurational methods



A Mountaineering Strategy to Excited States: Highly Accurate Energies and Benchmarks for Medium Sized Molecules

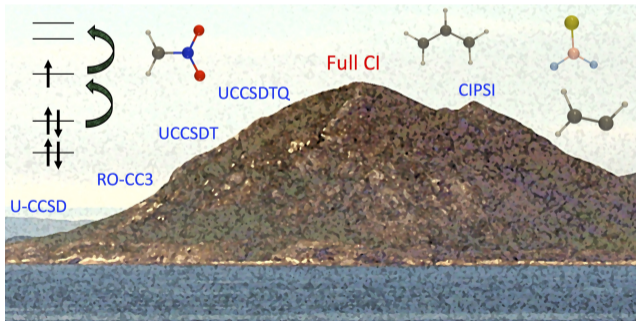
Pierre-François Loos,* Filippo Lipparini,* Martial Boggió-Pasqua, Anthony Scemama, and Denis Jacquemin*



- ▶ 238 VTEs (and oscillator strengths) with mostly singly-excited states and aug-cc-pVTZ basis
- ▶ 27 small- and medium-sized molecules with 4 to 6 non-H atoms
- ▶ rely mostly on **CCSDT** or CCSDTQ to define TBEs
- ▶ benchmark popular excited-state methods accounting for double and triple excitations
- ▶ recently improved TBEs with CC4 and CCSDTQ [JCP 154 (2021) 221103; JCTC 18 (2022) 4418]

Mountaineering Strategy to Excited States: Highly Accurate Energies and Benchmarks for Exotic Molecules and Radicals

Pierre-François Loos,* Anthony Scemama, Martial Boggio-Pasqua, and Denis Jacquemin*



- ▶ two subsets of excitations and oscillator strengths
- ▶ an “exotic” subset of 30 VTEs for closed-shell molecules containing F, Cl, P, and Si
- ▶ a “radical” subset of 51 doublet-doublet transitions in 24 small radicals
- ▶ total of 81 TBEs mostly obtained at the FCI/aug-cc-pVTZ level
- ▶ benchmark popular excited-state methods (U vs RO)

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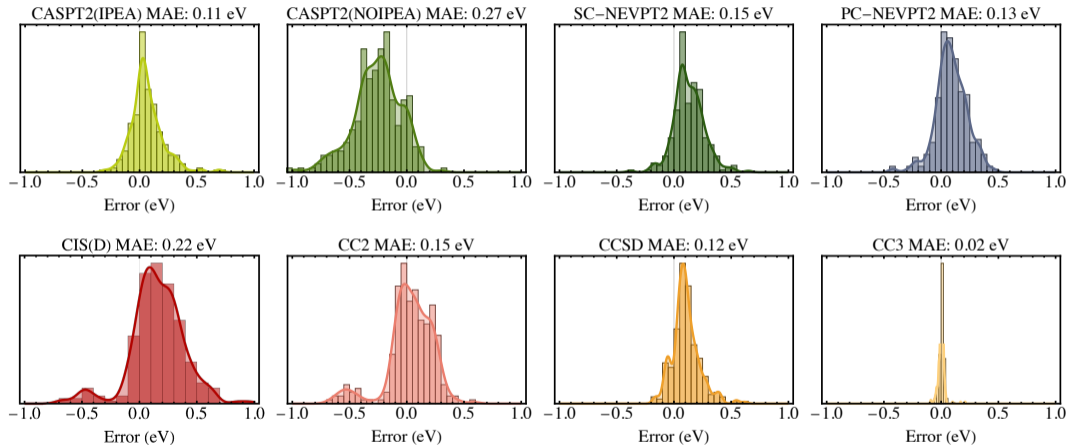


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QUESTDB: A database of highly accurate excitation energies for the electronic structure community

Mickaël Vэрil¹ | Anthony Scemama¹ | Michel Caffarel¹ | Filippo Lipparini² | Martial Boggio-Pasqua¹ | Denis Jacquemin³ | Pierre-François Loos¹

- ▶ 13 new systems composed by small molecules as well as **larger** molecules
- ▶ 80 new transitions the vast majority being of CCSDT quality
- ▶ benchmark popular excited-state methods **over the entire database**



JCTC 18 (2022) 2418; JCP 157 (2022) 014103

QUEST: a database of highly-accurate excitation energies

HOME DATASET SUBSETS REFERENCES

QUESTDB

Quantum Excited States DataBase

FCI/CBS
Selected

Small-size molecules

Medium-size molecules

AVTZ
AVDZ

NEVPT2

GROUND STATE

CASPT2

CCSDTQ

Full CI

CIPSI

UCCSDTQ

UCCSDT

A mountaineering strategy to excited states



Mika Vériel

Vériel et al. WIREs Comput. Mol. Sci. 11 (2021) e1517

https://lcpq.github.io/QUESTDB_website

- ▶ Head-Gordon's group: orbital-optimized DFT for double excitations [JCTC 16 (2020) 1699; JPCL 12 (2021) 4517] and TD-DFT benchmark [JCTC 18 (2022) 3460]
- ▶ Kaupp's group: assessment of hybrid functionals [JCP 155 (2021) 124108]
- ▶ Kallay's and Goerigk's groups: double hybrids [JCTC 15 (2019) 4735; JCTC 17 (2021) 927; JCTC 17 (2021) 5165; JCTC 17 (2021) 4211]
- ▶ Truhlar/Gagliardi's group: p-DFT [JCTC 18 (2022) 6065]
- ▶ Bartlett's group: Variants of EOM-CC for doubly-excited states [JCP 156 (2022) 201102; JPCA 127 (2023) 828; JCP 159 (2023) 094101]
- ▶ Neuscamman's group: QMC for doubly-excited states [JCP 153 (2022) 234105]
- ▶ Filippi/Scemama's groups: QMC for excited states [JCTC 15 (2019) 4889; JCTC 17 (2021) 3426; JCTC 18 (2022) 1089; JCTC 18 (2022) 6722]
- ▶ Gould's group: ensemble DFT [JPCL 13 (2022) 2452]

- ▶ Antoine Marie, Enzo Monino, Roberto Orlando, Yann Damour & Mika Véril
- ▶ Sara Giarrusso, Raúl Quintero-Monsebaiz & Fábris Kossoski
- ▶ Anthony Scemama
- ▶ Denis Jacquemin
- ▶ Martial Boggio-Pasqua
- ▶ Michel Caffarel



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