

State-Specific Coupled Cluster for Excited States

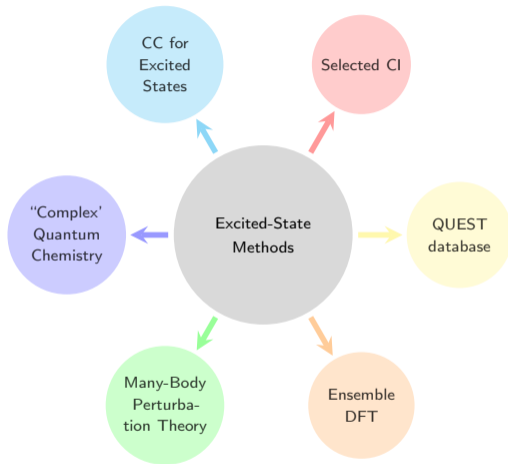
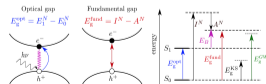
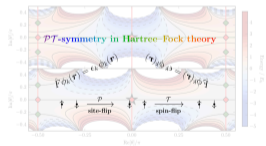
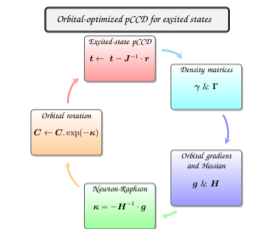
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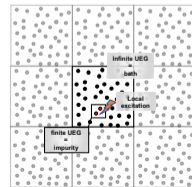
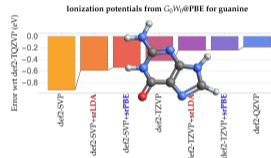
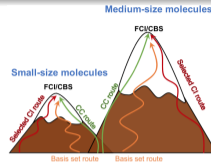
WATOC 2020 (Vancouver, BC)



General Overview of our Research Group



<https://lcpq.github.io/PTEROSOR/>





Fábri Kossoski
(Postdoc)



Antoine Marie
(PhD)

- Coupled-cluster (CC) wave function

$$|\Psi_{\text{CC}}\rangle = e^{\hat{T}} |\Psi_0\rangle \quad \text{where } \Psi_0 \text{ is a reference wave function} \quad (1)$$

- Exponential *ansatz*

$$e^{\hat{T}} = \hat{I} + \hat{T} + \frac{1}{2!} \hat{T}^2 + \frac{1}{3!} \hat{T}^3 + \dots \quad (2)$$

- Excitation operator

$$\hat{T} = \hat{T}_1 + \hat{T}_2 + \dots + \hat{T}_n \quad \text{where } n \text{ is the number of electrons} \quad (3)$$

$$\hat{T}_1 |\Psi_0\rangle = \sum_i \sum_a \underbrace{t_i^a}_{\text{amplitudes}} |\Psi_i^a\rangle \quad \hat{T}_2 |\Psi_0\rangle = \sum_{i < j} \sum_{a < b} t_{ij}^{ab} \underbrace{|\Psi_{ij}^{ab}\rangle}_{\text{excited determinants}} \quad (4)$$

- Schrödinger equation

$$\hat{H}|\Psi_{CC}\rangle = E_{CC}|\Psi_{CC}\rangle \Rightarrow \hat{H}e^{\hat{T}}|\Psi_0\rangle = E_{CC}e^{\hat{T}}|\Psi_0\rangle \Rightarrow \underbrace{e^{-\hat{T}}\hat{H}e^{\hat{T}}}_{\bar{H} = \text{similarity transform}}|\Psi_0\rangle = E_{CC}|\Psi_0\rangle \quad (5)$$

- Variational CC energy (**factorial complexity**)

$$E_{VCC} = \frac{\langle\Psi_{CC}|\hat{H}|\Psi_{CC}\rangle}{\langle\Psi_{CC}|\Psi_{CC}\rangle} = \frac{\langle\Psi_0(e^{\hat{T}})^\dagger|\hat{H}|e^{\hat{T}}\Psi_0\rangle}{\langle\Psi_0(e^{\hat{T}})^\dagger|e^{\hat{T}}\Psi_0\rangle} \geq E_{\text{exact}} \quad (6)$$

- (Traditional) projected CC energy (**polynomial complexity**)

$$E_{TCC} = \frac{\langle\Psi_0|\bar{H}|\Psi_0\rangle}{\langle\Psi_0|\Psi_0\rangle} = \frac{\langle\Psi_0e^{-\hat{T}}|\hat{H}|e^{\hat{T}}\Psi_0\rangle}{\langle\Psi_0e^{-\hat{T}}|e^{\hat{T}}\Psi_0\rangle} \quad (7)$$

- Unitary CC energy (**very expensive unless you have a quantum computer**)

$$E_{UCC} = \frac{\langle\Psi_0(e^{\hat{\tau}})^\dagger|\hat{H}|e^{\hat{\tau}}\Psi_0\rangle}{\langle\Psi_0(e^{\hat{\tau}})^\dagger|e^{\hat{\tau}}\Psi_0\rangle} = \frac{\langle\Psi_0e^{-\hat{\tau}}|\hat{H}|e^{\hat{\tau}}\Psi_0\rangle}{\langle\Psi_0|\Psi_0\rangle} \quad \text{where } \hat{\tau} = \hat{T} - \hat{T}^\dagger \text{ is anti-Hermitian} \quad (8)$$

- Amplitude equations for **single amplitudes**

$$\underbrace{r_i^a}_{\text{residual}} = \langle \Psi_i^a | \bar{H} | \Psi_0 \rangle = 0 \Rightarrow t_i^a \quad (9)$$

- Amplitude equations for **double amplitudes**

$$r_{ij}^{ab} = \langle \Psi_{ij}^{ab} | \bar{H} | \Psi_0 \rangle = 0 \Rightarrow t_{ij}^{ab} \quad (10)$$

- Amplitude equations for **k-tuple amplitudes**

$$r_{ij\dots}^{ab\dots} = \langle \Psi_{ij\dots}^{ab\dots} | \bar{H} | \Psi_0 \rangle = 0 \Rightarrow t_{ij\dots}^{ab\dots} \quad (11)$$

Traditional pair CCD (TpCCD)

- We set $\hat{T} = \hat{T}_2$ (CCD) and we restrict the excitation manifold to electron pairs (p)

$$\hat{T} |\Psi_0\rangle = \sum_{ia} t_{ii}^{a\bar{a}} |\Psi_{ii}^{a\bar{a}}\rangle = \sum_{ia} t_i^a |\Psi_{ii}^{a\bar{a}}\rangle \quad (12)$$

- TpCCD energy

$$E_{\text{TpCCD}} = \langle \Psi_0 | \hat{H} | \Psi_0 \rangle + \sum_{ia} t_i^a \langle ii | aa \rangle \quad (13)$$

- TpCCD residual

$$\begin{aligned} r_i^a = & \langle ii | aa \rangle + 2(f_a^a - f_i^i) t_i^a - 2 \sum_j \langle jj | aa \rangle t_j^a t_i^a - 2 \sum_b \langle ii | bb \rangle t_i^b t_i^a \\ & - 2(2 \langle ia | ia \rangle - \langle ia | ai \rangle) t_i^a + 2 \langle ii | aa \rangle t_i^a t_i^a \\ & + \sum_b \langle aa | bb \rangle t_i^b + \sum_j \langle ii | jj \rangle t_j^a + \sum_{jb} \langle jj | bb \rangle t_j^a t_i^b \end{aligned} \quad (14)$$

- Update amplitudes

$$t_i^a \leftarrow t_i^a - \frac{r_i^a}{2f_a^a - 2f_i^i} \quad \text{quasi-Newton algorithm} \quad (15)$$

Limacher et al. JCTC 9 (2013) 1394, Henderson et al. JCP 141 (2014) 244104, and many others.

How to target excited states at the CC level?

Equation-of-motion (EOM) formalism

$$\bar{H} = \begin{pmatrix} E_{CC} & \langle \Psi_0 | \bar{H} | \Psi_i^a \rangle & \langle \Psi_0 | \bar{H} | \Psi_{ij}^{ab} \rangle \\ 0 & \langle \Psi_i^a | \bar{H} | \Psi_i^a \rangle & \langle \Psi_i^a | \bar{H} | \Psi_{ij}^{ab} \rangle \\ 0 & \langle \Psi_{ij}^{ab} | \bar{H} | \Psi_i^a \rangle & \langle \Psi_{ij}^{ab} | \bar{H} | \Psi_{ij}^{ab} \rangle \end{pmatrix}$$

This is biased towards the ground state!!

Stanton & Bartlett, JCP 98 (1993) 7029

"State-specific" formalism

$$r_i^a = \langle \Psi_i^a | \bar{H} | \Psi_0 \rangle = 0$$
$$r_{ij}^{ab} = \langle \Psi_{ij}^{ab} | \bar{H} | \Psi_0 \rangle = 0$$

There is more than one solution!!

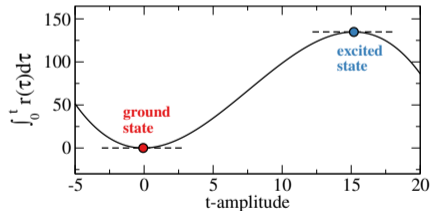
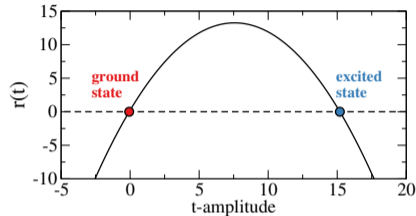
Piecuch & Kowalski, in Computational Chemistry: Reviews of Current Trends, Vol. 5 (2000) 1

We're looking for the roots of $r(t)$...

Residual and curvature

$$\begin{aligned}
 r(t) = & \langle hh|ll \rangle \\
 & + \left(2f_l^l - 2f_h^h - 4 \langle hl|hl \rangle + 2 \langle lh|hl \rangle + \langle ll|ll \rangle + \langle hh|hh \rangle \right) t \\
 & - \langle ll|hh \rangle t^2
 \end{aligned} \tag{16}$$

$$\begin{aligned}
 r'(t) = & 2f_l^l - 2f_h^h - 4 \langle hl|hl \rangle + 2 \langle lh|hl \rangle + \langle ll|ll \rangle + \langle hh|hh \rangle \\
 & - 2 \langle ll|hh \rangle t
 \end{aligned} \tag{17}$$



$$t_i^a \leftarrow t_i^a - \sum_{jb} (J^{-1})_{ia,jb} r_j^b \quad \text{Newton-Raphson algorithm} \quad (18)$$

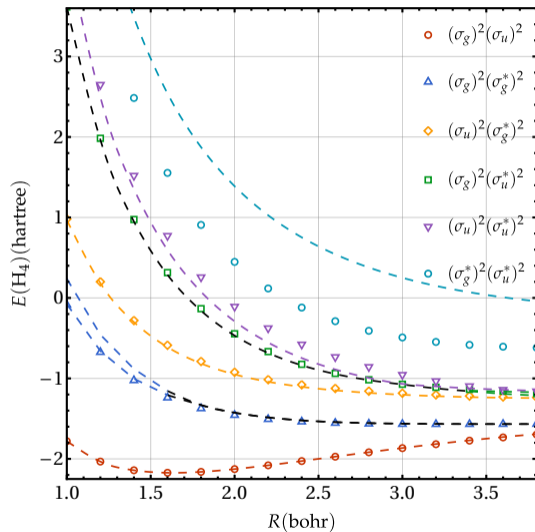
Elements of the exact Jacobian matrix for TpCCD

$$\begin{aligned}
 J_{ia,jb} = \frac{\partial r_i^a}{\partial t_j^b} = & \left[2f_a^a - 2f_i^i - 4 \langle ia|ia \rangle + 2 \langle ia|ai \rangle \right] \delta_{ij} \delta_{ab} \\
 & + \left[\langle aa|bb \rangle - \langle jj|aa \rangle t_i^a + (1 - 2\delta_{ab}) \sum_{k \neq i} \langle kk|bb \rangle t_k^a \right] \delta_{ij} \\
 & + \left[\langle ii|jj \rangle - \langle ii|bb \rangle t_i^a + (1 - 2\delta_{ij}) \sum_{c \neq a} \langle jj|cc \rangle t_i^c \right] \delta_{ab}.
 \end{aligned} \quad (19)$$

Kossoski et al. JCTC 17 (2021) 4756

NB: Same strategy works for VCC [Marie et al. JCP 155 (2021) 104105]

Stretching linear H₄/STO-6G: pCCD vs DOCI with HF ground-state reference



Orbital rotations via unitary transformation

$$\tilde{E}(\hat{T}, \hat{\kappa}) = \langle \Psi_0 | \underbrace{(\hat{I} + \hat{Z})}_{\text{de-excitation operator}} e^{-\hat{T}} e^{-\hat{\kappa}} \hat{H} e^{\hat{\kappa}} e^{\hat{T}} | \Psi_0 \rangle \quad (20)$$

NB: pCCD is not invariant wrt orbital rotations...

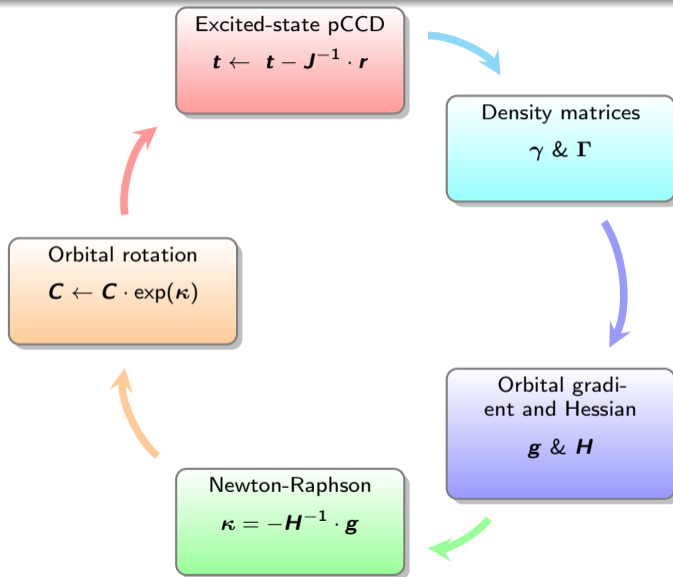
Updating the orbital coefficients

$$\tilde{E}(\kappa) \approx \tilde{E}(0) + \mathbf{g} \cdot \kappa + \frac{1}{2} \kappa^\dagger \cdot \mathbf{H} \cdot \kappa \Rightarrow \underbrace{\mathbf{C}}_{\text{new coefficients!}} \leftarrow \mathbf{C} \cdot e^\kappa \quad \text{with} \quad \kappa = -\mathbf{H}^{-1} \cdot \mathbf{g} \quad (21)$$

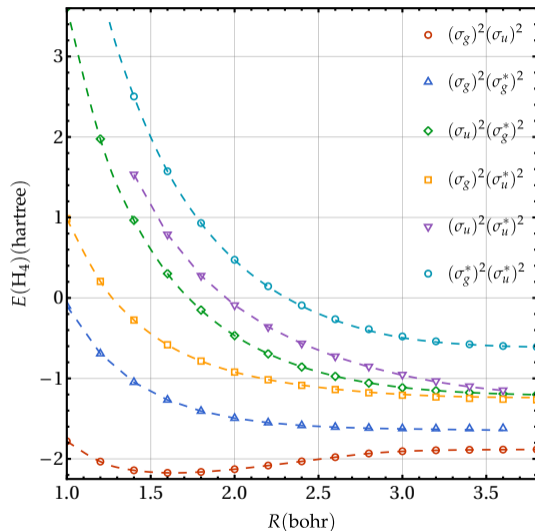
$$\text{Density matrices } \gamma \text{ \& } \Gamma \Rightarrow \underbrace{g_{pq}}_{\text{gradient}} = \left. \frac{\partial \tilde{E}(\kappa)}{\partial \kappa_{pq}} \right|_{\kappa=0} \quad \underbrace{H_{pq,rs}}_{\text{Hessian}} = \left. \frac{\partial^2 \tilde{E}(\kappa)}{\partial \kappa_{pq} \partial \kappa_{rs}} \right|_{\kappa=0} \quad (22)$$

Henderson et al. JCP 141 (2014) 244104

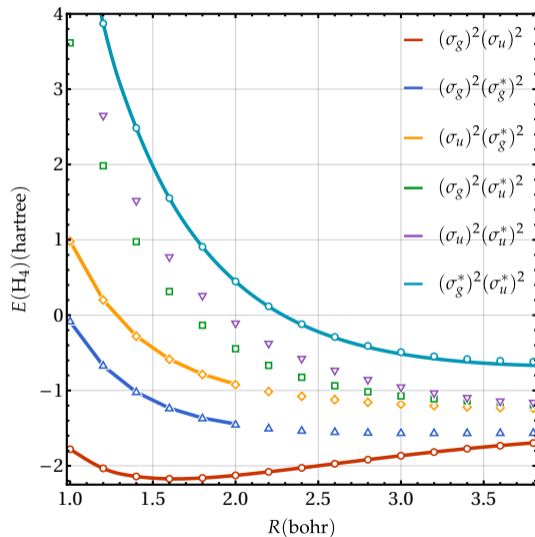
State-specific orbital-optimized TpCCD (oo-TpCCD) for excited states



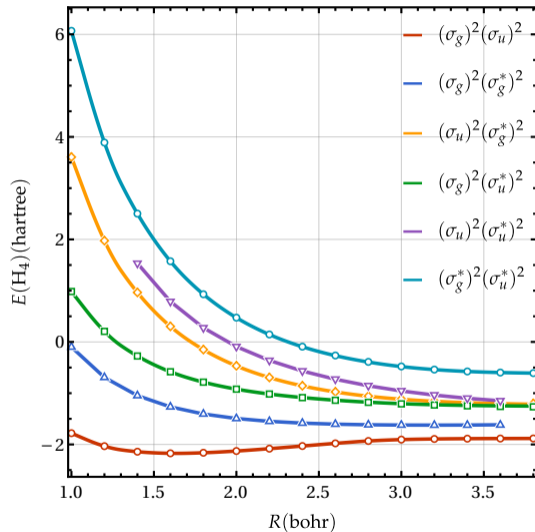
Stretching linear H₄/STO-6G: TpCCD vs DOCI with state-specific TpCCD reference



Stretching linear H₄/STO-6G: VpCCD vs DOCI with HF ground-state reference



Stretching linear H₄/STO-6G: VpCCD vs DOCI with state-specific VpCCD reference



Doubly-excited states for molecules (6-31+G*)...

molecule	method	ΔE (eV)	$\Delta\Delta E$ (eV)
BH	Δ_{oo} -TpCCD	7.35	+0.24
	FCI	7.11	0
	EOM-CCSDTQ	7.12	+0.01
	EOM-CCSDT	7.15	+0.04
	CC3	7.30	+0.19
HNO	Δ_{oo} -TpCCD	4.49	-0.02
	FCI ¹	4.51	0
	EOM-CCSDTQ ¹	4.54	+0.03
	EOM-CCSDT ¹	4.81	+0.30
	CC3 ¹	5.28	+0.77
H ₃ C-NO	Δ_{oo} -TpCCD	4.66	-0.20
	FCI ¹	4.86	0
	EOM-CCSDT ¹	5.26	+0.40
	CC3 ¹	5.73	+0.87
H ₂ C=O	Δ_{oo} -TpCCD	11.26	+0.40
	FCI ¹	10.86	0
	EOM-CCSDTQ ¹	10.87	+0.01
	EOM-CCSDT ¹	11.10	+0.24
	CC3 ¹	11.49	+0.63

¹Results and geometries from JCTC 15 (2019) 1939

Quantum Package 2.0: <https://github.com/QuantumPackage/qp2>

The image is a composite of three parts illustrating the Quantum Package 2.0 environment:

- Terminal Window (Left):** Shows the command-line interface. It includes the `qp_plugins` command to list and manage plugins, the `qpsh` shell prompt, and a sequence of commands to create and run a calculation for ethanol (`ethanol.xyz`), including setting up the file, running the program, and converting the output.
- 3D Box (Center):** A 3D rendering of a cardboard box labeled "QUANTUM PACKAGE 2.0" with a stylized atomic symbol on its side.
- Output Window (Right):** Displays a molecular orbital visualization of an ethanol molecule (CCO) with a 3D ball-and-stick model. Below the model is a table of numerical data, likely representing orbital energies or coefficients.

*"Quantum Package 2.0: An Open-Source Determinant-Driven Suite of Programs",
Garniron et al., JCTC 15 (2019) 3591*

Fábris' repo: https://github.com/kossoski/qp_plugins_kossoski

- Fábris Kossoski
[Kossoski et al. JCTC 17 (2021) 4756]
- Antoine Marie
[Marie et al. JCP 155 (2021) 104105]
- Raul Quintero (Poster #739 on Fock-space CC)
- Enzo Monino (Poster #524 on GW methods)
[Monino & Loos JCP 156 (2022) 231101]
- Yann Damour (SCI methods with optimized orbitals)
[Damour et al. JCP 155 (2021) 134104]



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<https://lcpq.github.io/PTEROSOR>