

GW, BSE, eDFT, RPA, and piña colada

Pierre-François LOOS

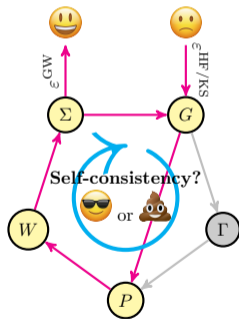
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Gori-Giorgi Symposium



Section 1

Many-body perturbation theory



arXiv:2002.04514 [physics.chem-ph]

Pros and Cons of the Bethe-Salpeter Formalism for Ground-State Energies

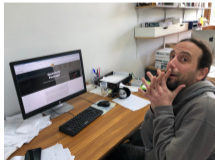
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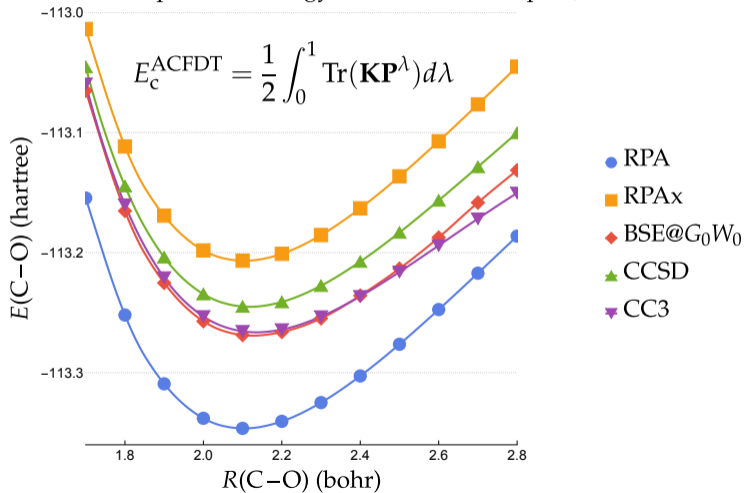
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Bethe-Salpeter for ground-state energies

Ground-state potential energy surface of CO/cc-pVQZ



BSE correlation energy via ACFDT:

$$E_c^{\text{BSE}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K}\mathbf{P}^\lambda) d\lambda$$

Correlation part of the two-electron density matrix:

$$\mathbf{P}^\lambda = \begin{pmatrix} \mathbf{Y}^\lambda (\mathbf{Y}^\lambda)^\top & \mathbf{Y}^\lambda (\mathbf{X}^\lambda)^\top \\ \mathbf{X}^\lambda (\mathbf{Y}^\lambda)^\top & \mathbf{X}^\lambda (\mathbf{X}^\lambda)^\top \end{pmatrix} - \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{1} \end{pmatrix}$$

“Casida”-like equations:

$$\begin{pmatrix} \mathbf{A}^\lambda & \mathbf{B}^\lambda \\ -\mathbf{B}^\lambda & -\mathbf{A}^\lambda \end{pmatrix} \begin{pmatrix} \mathbf{X}_m^\lambda \\ \mathbf{Y}_m^\lambda \end{pmatrix} = \Omega_m^\lambda \begin{pmatrix} \mathbf{X}_m^\lambda \\ \mathbf{Y}_m^\lambda \end{pmatrix}$$

with

$$A_{ia,jb}^{\lambda,\text{BSE}} = \delta_{ij}\delta_{ab}(\epsilon_a^{\text{GW}} - \epsilon_i^{\text{GW}}) + \lambda \left[2(ia|jb) - W_{ij,ab}^\lambda \right]$$

$$B_{ia,jb}^{\lambda,\text{BSE}} = \lambda \left[2(ia|bj) - W_{ib,aj}^\lambda \right]$$

Screened Coulomb operator:

$$W_{ij,ab}^{\lambda}(\omega) = (ij|ab) + 2 \sum_m^{OV} [ij|m]^{\lambda} [ab|m]^{\lambda} \left(\frac{1}{\omega - \Omega_m^{\lambda, \text{RPA}} + i\eta} - \frac{1}{\omega + \Omega_m^{\lambda, \text{RPA}} - i\eta} \right)$$

Screened two-electron integrals:

$$[pq|m]^{\lambda} = \sum_i^O \sum_a^V (pq|ia) (\mathbf{X}_m^{\lambda} + \mathbf{Y}_m^{\lambda})_{ia}$$

Bethe-Salpeter for ground-state energies

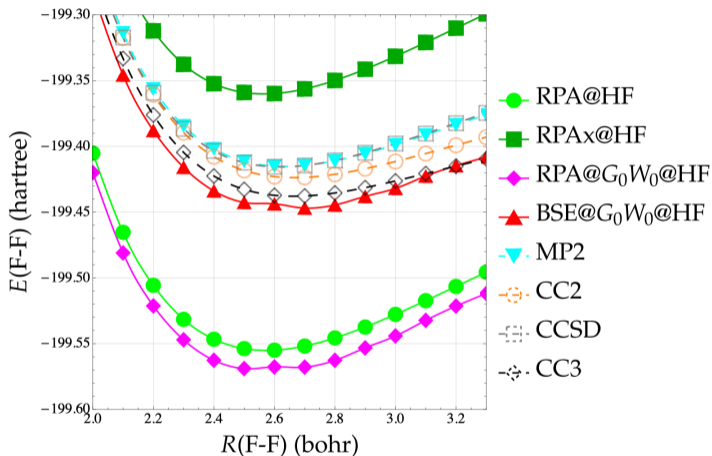


FIG. 4. Ground-state PES of F_2 around its equilibrium geometry obtained at various levels of theory with the cc-pVQZ basis set.

Green Functions and Self-Consistency: Insights From the Spherium Model

Pierre-François Loos,^{*,†} Pina Romaniello,^{‡,¶} and J. A. Berger^{†,¶}

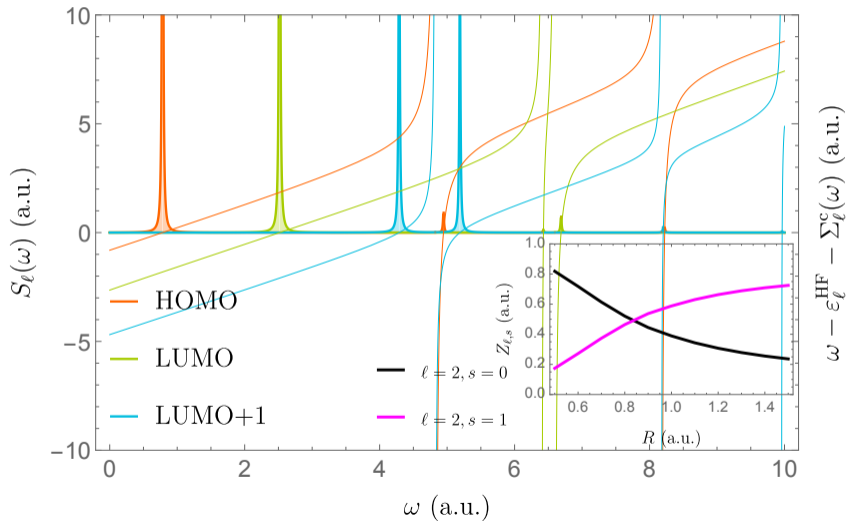
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The elephant in the room of GW



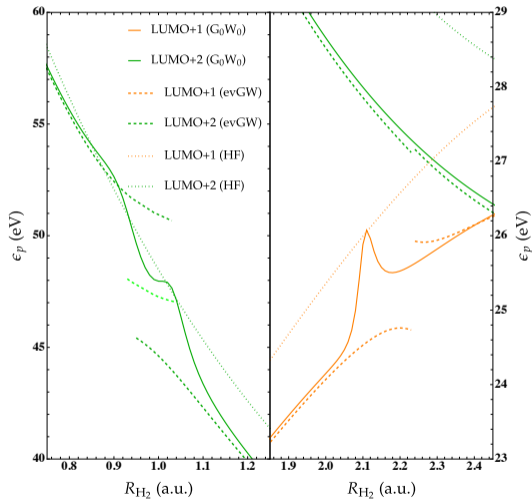
Unphysical Discontinuities in GW Methods

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The elephant in the room of GW



(Linearized) quasiparticle equation

$$\epsilon_p^{G_0 W_0} = \epsilon_p^{\text{HF}} + Z_p(\epsilon_p^{\text{HF}}) \text{Re}[\Sigma_p^c(\epsilon_p^{\text{HF}})]$$

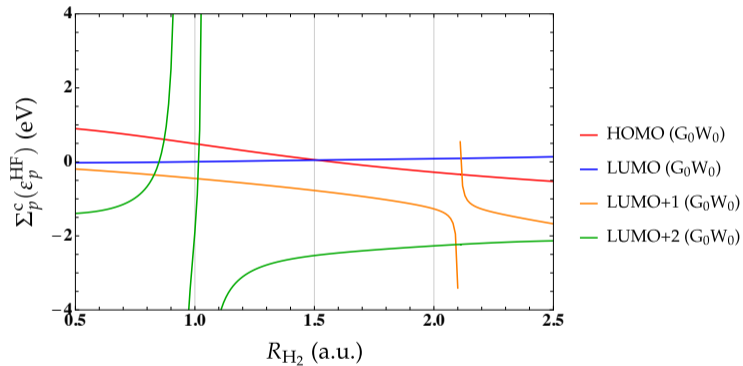
Correlation part of the self-energy:

$$\Sigma_p^c(\omega) = 2 \sum_{im} \frac{[pi|m]^2}{\omega - \epsilon_i^{\text{HF}} + \Omega_m^{\text{RPA}} - i\eta} + 2 \sum_{am} \frac{[pa|m]^2}{\omega - \epsilon_a^{\text{HF}} - \Omega_m^{\text{RPA}} + i\eta}$$

Renormalization factor (or spectral weight):

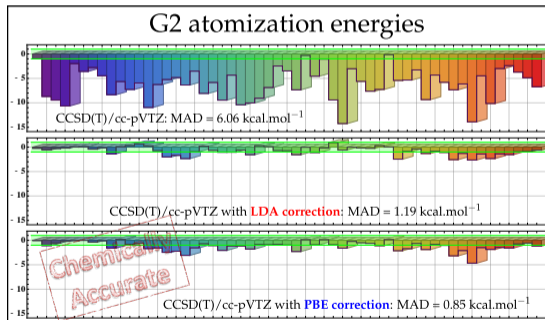
$$Z_p(\omega) = \left[1 - \frac{\partial \text{Re}[\Sigma_p^c(\omega)]}{\partial \omega} \right]^{-1}$$

The elephant in the room of GW



Section 2

Basis set incompleteness correction



A Density-Based Basis-Set Correction for Wave Function Theory

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and Emmanuel Giner^{*,‡}

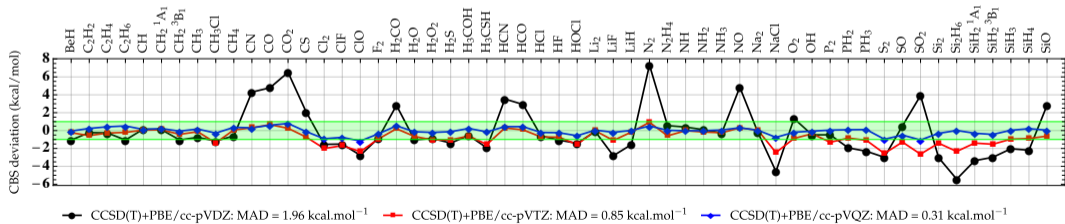
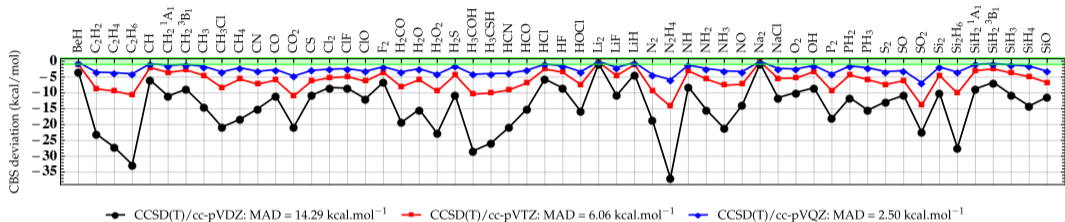
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Density-based basis set incompleteness error



Chemically accurate excitation energies with small basis sets

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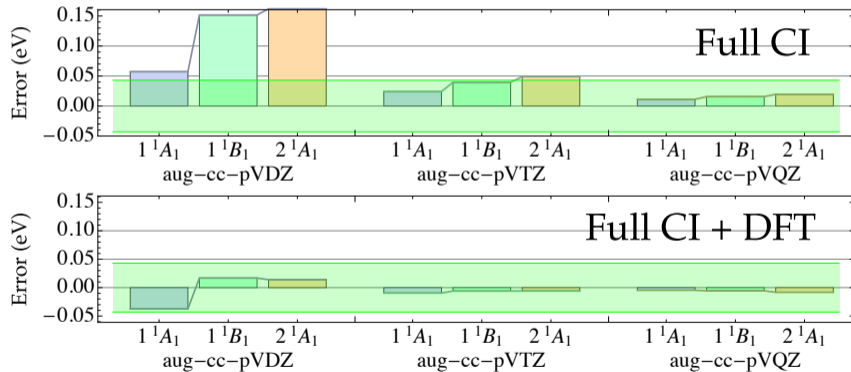
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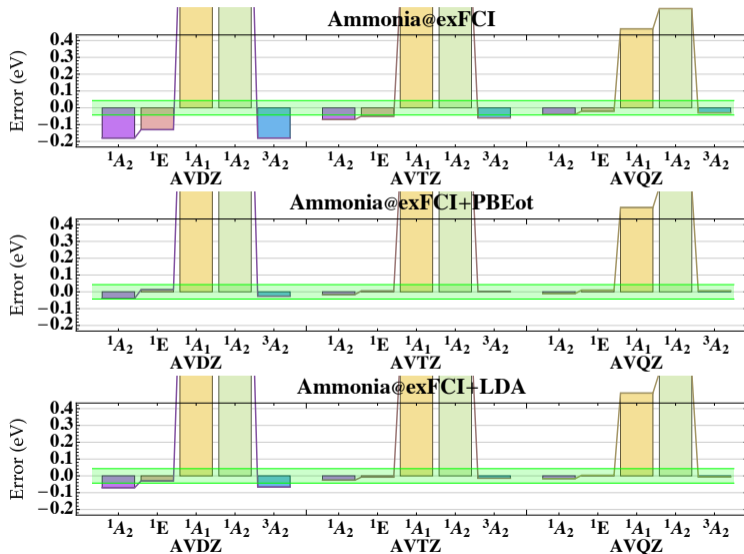


Density-based basis set incompleteness error

Adiabatic energies of methylene



Density-based basis set incompleteness error



Density-Based Basis-Set Incompleteness Correction for *GW* Methods

Pierre-François Loos,* Barthélemy Pradines, Anthony Scemama, Emmanuel Giner, and Julien Toulouse*



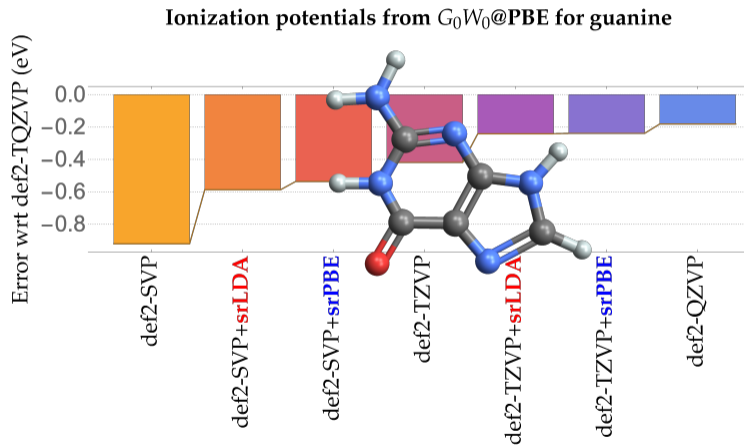
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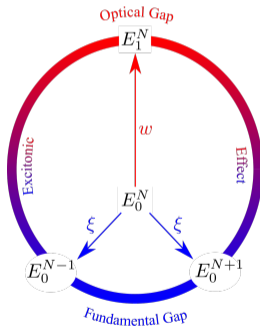


Density-based basis set incompleteness error



Section 3

Density-functional theory for ensembles



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COMING SOON



Gross-Oliveira-Kohn (GOK) DFT in a three-state ensemble

Ensemble energy:

$$E^{\mathbf{w}} = (1 - w_1 - w_2)E^{(0)} + w_1E^{(1)} + w_2E^{(2)}$$

Excitation energies:

$$\frac{\partial E^{\mathbf{w}}}{\partial w_1} = E^{(1)} - E^{(0)} = \Omega^{(1)} \quad \frac{\partial E^{\mathbf{w}}}{\partial w_2} = E^{(2)} - E^{(0)} = \Omega^{(2)}$$

Ensemble energy in practice:

$$E^{\mathbf{w}} = \min_n \left\{ F^{\mathbf{w}}[n] + \int v_{\text{ext}}(\mathbf{r})n(\mathbf{r})d\mathbf{r} \right\} \quad F^{\mathbf{w}}[n] = T_s^{\mathbf{w}}[n] + E_{\text{Hxc}}^{\mathbf{w}}[n]$$

Derivative discontinuity:

$$\frac{\partial E^{\mathbf{w}}}{\partial w_I} = \mathcal{E}_I^{\mathbf{w}} - \mathcal{E}_0^{\mathbf{w}} + \left. \frac{\partial E_{\text{xc}}^{\mathbf{w}}[n]}{\partial w_I} \right|_{n=n^{\mathbf{w}}(\mathbf{r})} \quad E_{\text{xc}}^{\mathbf{w}}[n] = \int \epsilon_{\text{xc}}^{\mathbf{w}}(n(\mathbf{r}))n(\mathbf{r})d\mathbf{r}$$

Construction of a weight-dependent LDA functional

Three-state ensemble exchange-correlation functional:

$$\tilde{\epsilon}_{\text{xc}}^{w_1, w_2}(n) = (1 - w_1 - w_2)\epsilon_{\text{xc}}^{(0)}(n) + w_1\epsilon_{\text{xc}}^{(1)}(n) + w_2\epsilon_{\text{xc}}^{(2)}(n)$$

LDA-centered functionals:

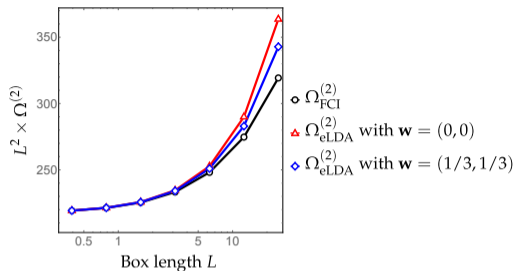
$$\bar{\epsilon}_{\text{xc}}^{(l)}(n) = \epsilon_{\text{xc}}^{(l)}(n) + \epsilon_{\text{xc}}^{\text{LDA}}(n) - \epsilon_{\text{xc}}^{(0)}(n)$$

$$\tilde{\epsilon}_{\text{xc}}^{w_1, w_2}(n) \rightarrow \epsilon_{\text{xc}}^{w_1, w_2}(n) = (1 - w_1 - w_2)\bar{\epsilon}_{\text{xc}}^{(0)}(n) + w_1\bar{\epsilon}_{\text{xc}}^{(1)}(n) + w_2\bar{\epsilon}_{\text{xc}}^{(2)}(n)$$

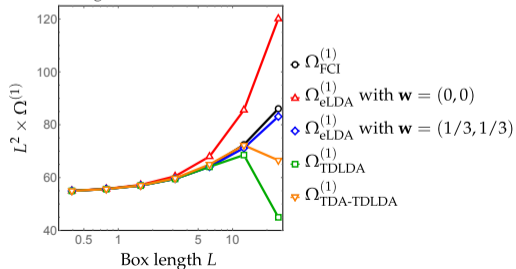
Weight-dependent LDA functional for ensembles “eLDA”:

$$\epsilon_{\text{xc}}^{w_1, w_2}(n) = \epsilon_{\text{xc}}^{\text{LDA}}(n) + w_1 \left[\epsilon_{\text{xc}}^{(1)}(n) - \epsilon_{\text{xc}}^{(0)}(n) \right] + w_2 \left[\epsilon_{\text{xc}}^{(2)}(n) - \epsilon_{\text{xc}}^{(0)}(n) \right]$$

Double excitation for $N = 5$



Single excitation for $N = 5$





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